

INTRODUCTION TO ROBOTICS

(Kinematics, Dynamics, and Design)

SESSION # 16: MANIPULATOR DYNAMICS

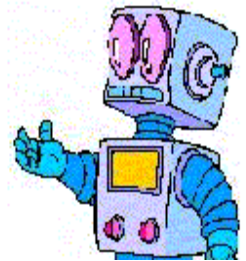
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Manipulator Dynamics

- So far we have only studied motion of manipulators without regard to forces causing the motion.
- Let us now derive the equations of motion for manipulator arms. In dynamics, we generally consider the following issues:
 - ❖ **Forward Dynamics:** Computing the resulting motion of the manipulator arm ($\theta, \dot{\theta}, \ddot{\theta}$) under the application of a set joint torques (τ). This is useful for simulation of the arm.
 - ❖ **Inverse Dynamics:** Computing the vector of joint torques (τ) for the given joint motion trajectory ($\theta, \dot{\theta}, \ddot{\theta}$). This is useful for controlling of the arm.



Manipulator Dynamics

Robotic Arms Dynamic Formulation History

Lagrangian Dynamics

Newton-Euler Dynamics

Kane Dynamics

Uicker/Kahn

Waters

Hollerbach

Kane/Levinson

(4x4) Matrices

Backward
Recursion
(4x4) Matrices

Forward
Recursion

Euler's Parameters
and
Relative Coordinates

(4x4) Matrices

(3x3) Matrices



Manipulator Dynamics

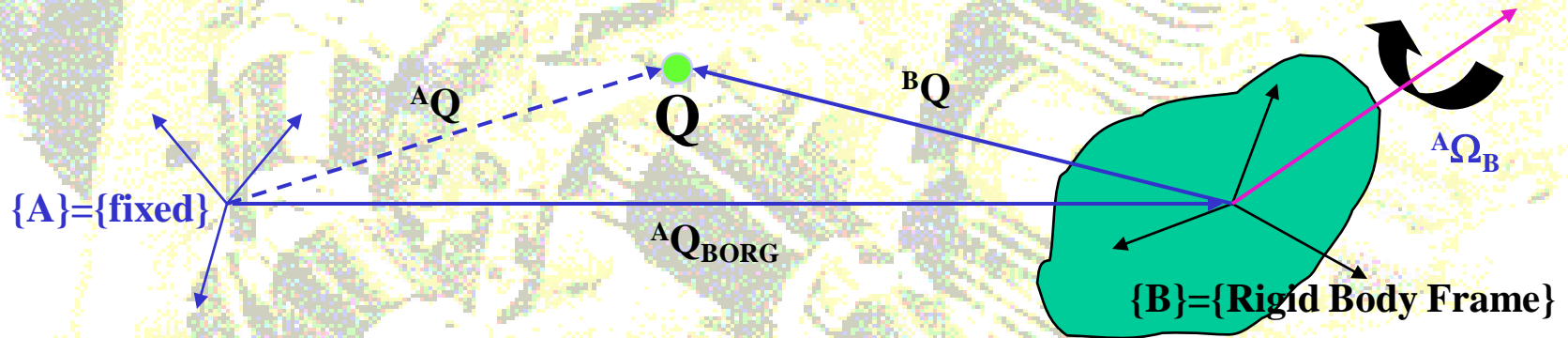
Author	Method	Multiplications	Additions
Uicker/Kahn (Lagrangian Dyn.)	(4 × 4) Matrices	66,271	51,548
Waters (Lagrangian Dyn.)	(4 × 4) Backward Recursion	7,051	5,652
Hollerbach (Lagrangian Dyn.)	(4 × 4) Forward Recursion	4,388	3,586
Hollerbach (Lagrangian Dyn.)	(3 × 3) Forward Recursion	2,195	1,719
Newton-Euler	Recursive	852	738
Kane/Levinson	Kane Dynamics	646	394
Raibert/Horn	Configuration Space Method (CSM)	468	264
Yang/Tzeng	Dyn. Simplification by Design	72	34 + 4 Trig. Functions.



Manipulator Dynamics

➤ Linear Accelerations of Rigid Bodies:

Consider a point “Q” in space, and describe its kinematics in two frames {A} and {B}.



From Chapter-5 we have:

$${}^A Q \equiv {}^A Q_{BORG} + {}^A R^B Q$$

$${}^A V_Q = {}^A V_{BORG} + {}^A \Omega_B \times {}^A R^B Q + {}^A R^B V_Q$$

Differentiating the velocity equation with respect to the time we have:



Manipulator Dynamics

➤ Linear Accelerations of Rigid Bodies:

Noting that: ${}^A\dot{R} \equiv {}^A\Omega_B \times {}^A R$

$${}^A\dot{V}_Q = {}^A\dot{V}_{BORG} + {}^A R^B \dot{V}_Q + 2 {}^A\Omega_B \times {}^A R^B V_Q + \\ + {}^A\dot{\Omega}_B \times {}^A R^B Q + {}^A\Omega_B \times ({}^A\Omega_B \times {}^A R^B Q)$$

If ${}^B Q$ is constant (on the R.B.), then: ${}^B V_Q = {}^B \dot{V}_Q = 0$

$${}^A\dot{V}_Q = {}^A\dot{V}_{BORG} + {}^A\dot{\Omega}_B \times {}^A R^B Q + {}^A\Omega_B \times ({}^A\Omega_B \times {}^A R^B Q)$$



Manipulator Dynamics

➤ Angular Acceleration of Rigid Bodies:

Consider:

- Frame **{B}** rotating relative to **{A}** with: ${}^A\Omega_B$

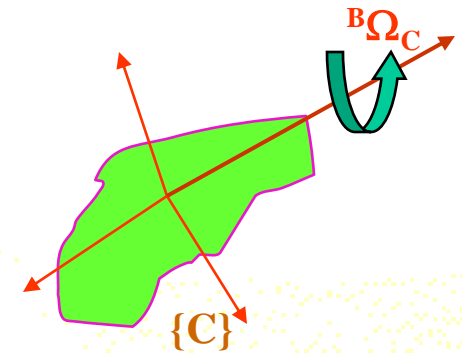
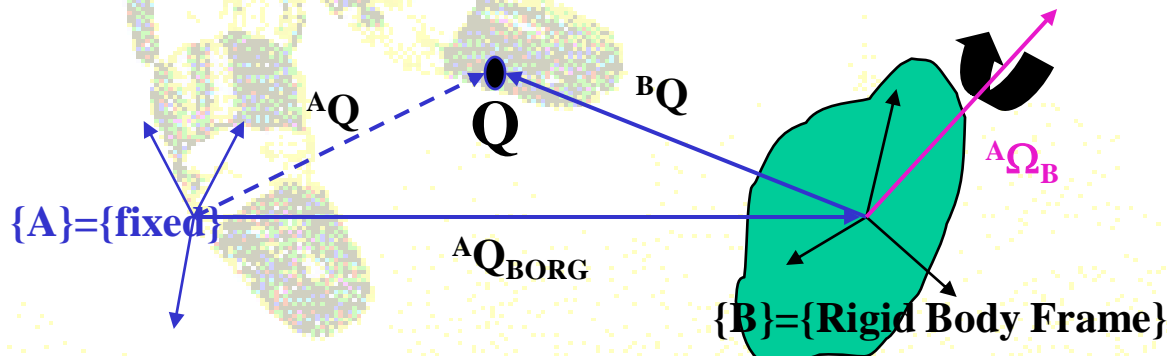
- Frames **{C}** rotating relative to **{B}** with: ${}^B\Omega_C$

Then:

$${}^A\Omega_C \equiv {}^A\Omega_B + {}^A R^B \Omega_C$$

Sum the vectors in frame {A}

$${}^A\dot{\Omega}_C = {}^A\dot{\Omega}_B + {}^A R^B \dot{\Omega}_C + {}^A\Omega_B \times {}^A R^B \Omega_C$$



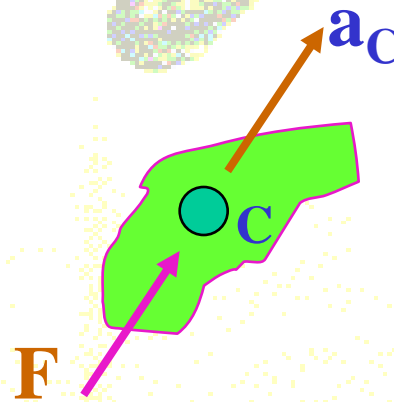
Manipulator Dynamics

➤ Newtonian Mechanics:

For a Rigid Body whose center of mass is accelerating with “ a_C ”, the Force “ F ” acting at the mass center is given by:

The Newton’s Law of Motion:

$$F = \sum f_i = \dot{P} = m\dot{v}_C = ma_C = (\text{Time rate of change of momentum})$$



Manipulator Dynamics

➤ Newtonian Mechanics:

For a Rigid Body rotating with an angular velocity “ ω ”, and an angular accelerating “ α ”, the Moment “ N ” which must be acting on the body to cause this motion, is given by:

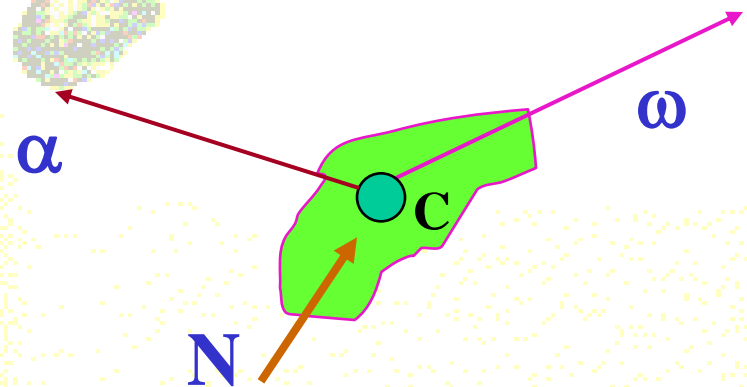
The Euler’s Equation:

$$N = {}^C I \cdot \alpha + \omega \times ({}^C I \cdot \omega)$$

(The rotational analogy of the Newton’s 2nd law comes from the Principle of Moment of Momentum)

where:

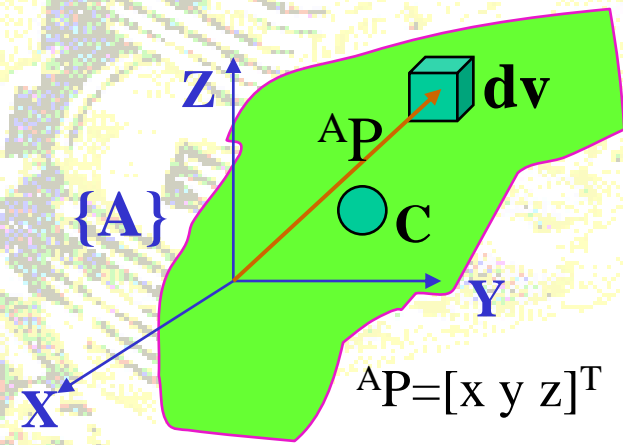
${}^C I$ = Inertia Tensor of the R.B.
written in frame {C}



Manipulator Dynamics

- **Mass Distribution:** The **Inertia Tensor** of an object describes the object's mass distribution (a generalization of the scalar moment of inertia). Relative to a frame $\{A\}$ is expressed as:

$${}^A I = \begin{bmatrix} I_{xx} & -I_{xy} & -I_{xz} \\ -I_{xy} & I_{yy} & -I_{yz} \\ -I_{xz} & -I_{yz} & I_{zz} \end{bmatrix}$$



where:

$$\begin{aligned} I_{xx} &= \iiint_v (y^2 + z^2) \rho dv; & I_{xy} &= \iiint_v xy \rho dv \\ I_{yy} &= \iiint_v (x^2 + z^2) \rho dv; & I_{xz} &= \iiint_v xz \rho dv \\ I_{zz} &= \iiint_v (x^2 + y^2) \rho dv; & I_{yz} &= \iiint_v yz \rho dv \end{aligned}$$



Manipulator Dynamics

➤ Iterative Newton-Euler Dynamic Formulation:

Let us now study the problem of computing the vector of joint torques (τ) for the given joint motion trajectory ($\theta, \dot{\theta}, \ddot{\theta}$). (The **Inverse Dynamics** problem useful for controlling of the arm).

❖ Outward Iterations to Compute Velocities and Accelerations:

To study dynamics from Newton & Euler equations, it is obvious that we need propagation equations for “ \dot{v} & $\dot{\omega}$ ”.

From Chapter-5, the angular velocity equation for every instant is:

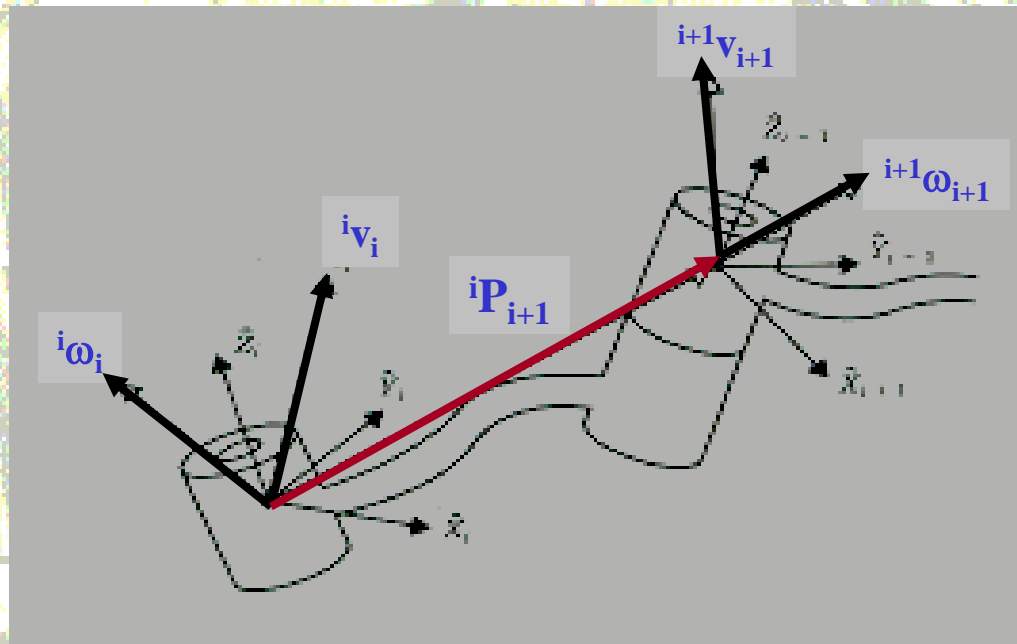
$${}^{i+1}\omega_{i+1} = {}^{i+1}R^i \omega_i + \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1}$$

Differentiating with respect to time we have:



Manipulator Dynamics

$${}^{i+1}\dot{\omega}_{i+1} = {}^{i+1}R^i \dot{\omega}_i + {}^{i+1}R^i \omega_i \times \dot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1} + \ddot{\theta}_{i+1} {}^{i+1}\hat{Z}_{i+1}$$



Where:

$${}^{i+1}\dot{\hat{Z}}_{i+1} = {}^{i+1}R^i \omega_i \times {}^{i+1}\hat{Z}_{i+1}$$

$${}^{i+1}\dot{R} = {}^{i+1}\omega_i \times {}^{i+1}R \Rightarrow {}^{i+1}\dot{R}^i \omega_i = {}^{i+1}\omega_i \times {}^{i+1}R^i \omega_i = {}^{i+1}\omega_i \times {}^{i+1}\omega_i = 0$$



Manipulator Dynamics

Also from Chapter-5, the linear velocity equation for every instant is:

$${}^{i+1}v_{i+1} = {}^i R({}^i v_i + {}^i \omega_i \times {}^i P_{i+1})$$

Differentiating with respect to time we have:

$${}^{i+1}\dot{v}_{i+1} = {}^i R({}^i \dot{v}_i + {}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times {}^i \dot{P}_{i+1}) \Rightarrow$$

$${}^{i+1}\dot{v}_{i+1} = {}^i R({}^i \dot{v}_i + {}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{i+1}))$$

Since at every instant:

$${}^i R = \text{constant} \Rightarrow {}^i \dot{R} = 0$$



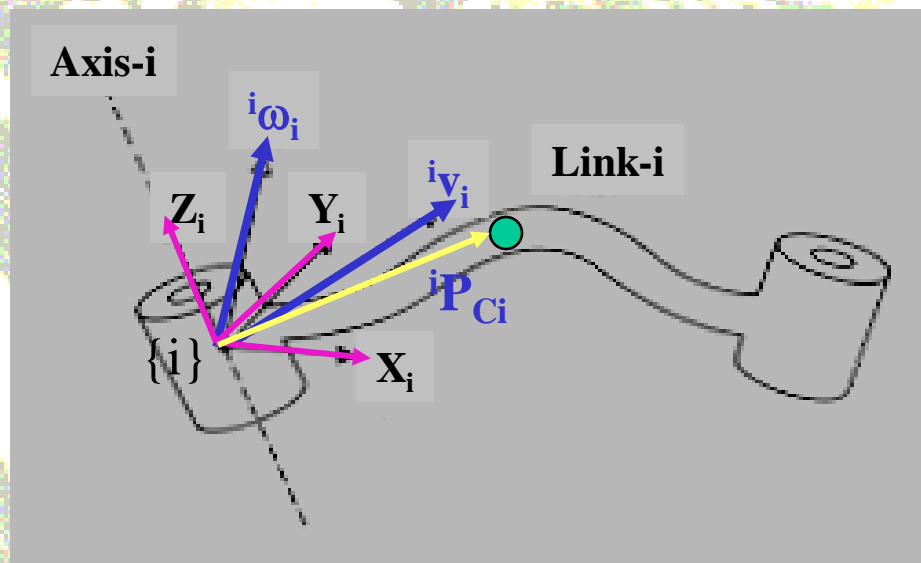
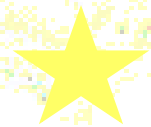
Manipulator Dynamics

To find the linear acceleration of the center of mass, we have:

$${}^i v_{C_i} = ({}^i v_i + {}^i \omega_i \times {}^i P_{C_i})$$

Differentiating with respect to time we have:

$${}^i \dot{v}_{C_i} = ({}^i \dot{v}_i + {}^i \dot{\omega}_i \times {}^i P_{C_i} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{C_i}))$$



Manipulator Dynamics

Having computed all acceleration equations, we shall now apply the **Newton-Euler Equations** as follows:

First compute the **Inertial Force and Torque** acting at the mass center of each link;

$$F_i = m\dot{v}_{C_i} = ma_{C_i}$$

$$N_i = {}^{C_i}I\dot{\omega}_i + \omega_i \times {}^{C_i}I\omega_i$$

$C_i I$ = Inertia Tensor of the link- i written in frame $\{C_i\}$ with it's origin at the mass center, and having the same orientation as frame $\{i\}$.

Then, perform **Inward Iterations** to compute forces and torques;

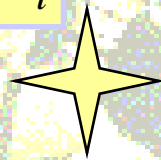


Manipulator Dynamics

❖ Inward Iterations to Compute Forces and Torques:

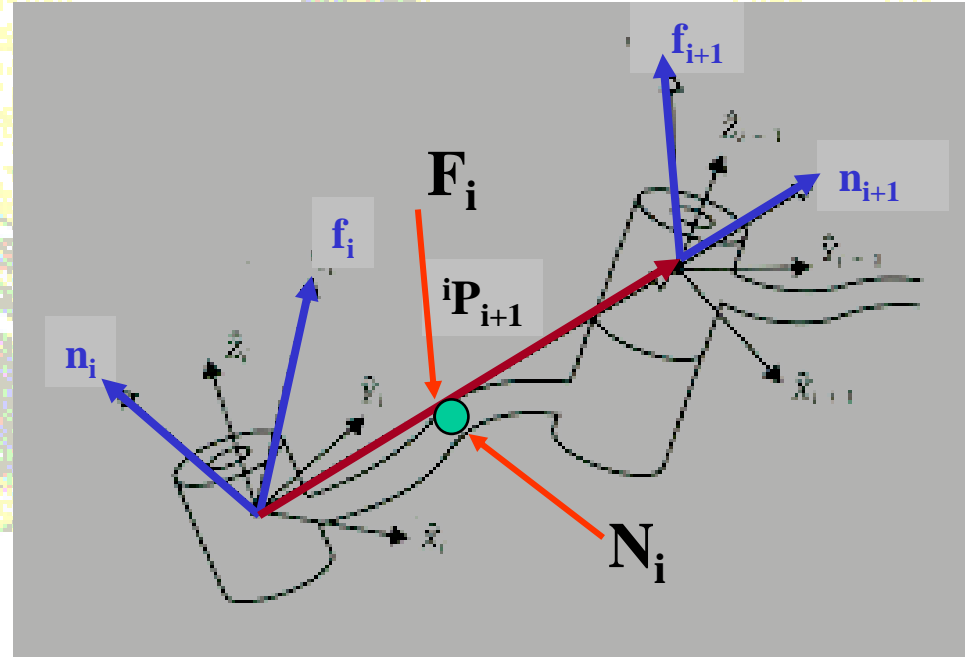
Write the force balance on link-i:

$${}^i f_i = {}^i R^{i+1} f_{i+1} + {}^i F_i$$



Write the moment balance about the origin of link frame-i:

$${}^i n_i = {}^i N_i + {}^i P_{i+1} \times {}^i R^{i+1} f_{i+1} + {}^i P_{Ci} \times {}^i F_i$$



Note: The required joint torques are found by taking the Z-component of the torque applied by one link on it's neighbor.

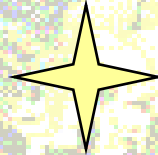


Manipulator Dynamics

❖ Inward Iterations to Compute Forces and Torques:

Therefore, for **Revolute Joints** we have:

$$\tau_i = {}^i n_i^T \hat{Z}_i$$



Therefore, for **Prismatic Joints** we have:

$$\tau_i = f_i^T \hat{Z}_i$$

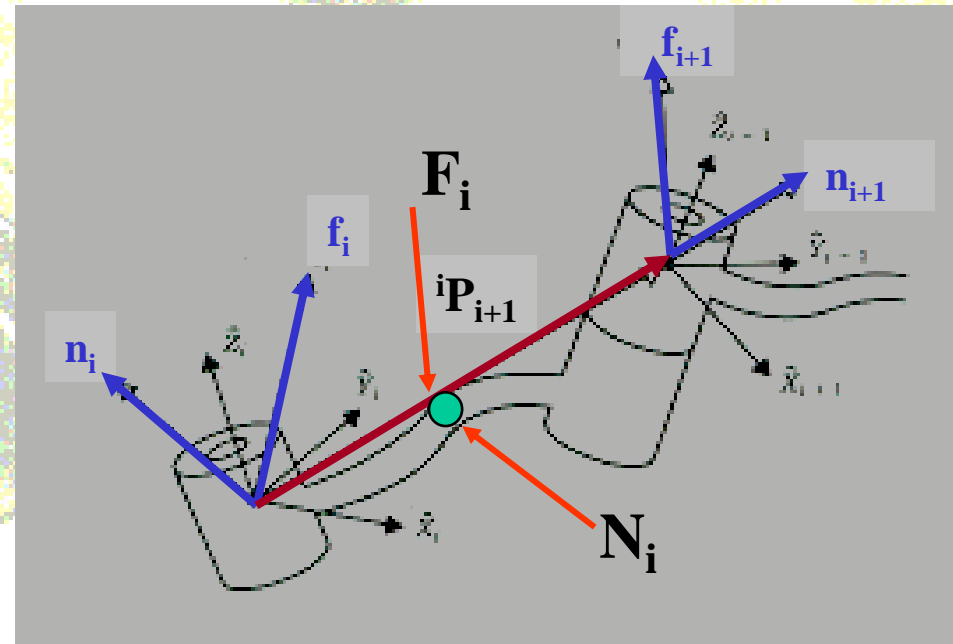


Note: For a robot moving in free space, we may have:

$${}^{N+1} f_{N+1} = {}^{N+1} n_{N+1} = 0$$

where as for a robot being in contact with the environment, we may have:

$${}^{N+1} f_{N+1} \neq {}^{N+1} n_{N+1} \neq 0$$



Manipulator Dynamics

➤ Iterative Newton-Euler Dynamic Algorithm:

First: Compute link velocities and accelerations iteratively from link-1 to link-n, and apply the Newton-Euler equations to each link.

Second: Compute the forces and torques of interaction recursively from link-n back to link-1.

Outward iterations: $i : 0 \rightarrow 5$

$${}^{i+1}\omega_{i+1} = {}^i R {}^i \omega_i + \dot{\theta}_{i+1} {}^{i+1} \hat{Z}_{i+1}, \quad (6.45)$$

$${}^{i+1}\dot{\omega}_{i+1} = {}^i R {}^i \dot{\omega}_i + {}^i R {}^i \omega_i \times \dot{\theta}_{i+1} {}^{i+1} \hat{Z}_{i+1} + \ddot{\theta}_{i+1} {}^{i+1} \hat{Z}_{i+1}, \quad (6.46)$$

$${}^{i+1}\dot{v}_{i+1} = {}^i R ({}^i \dot{\omega}_i \times {}^i P_{i+1} + {}^i \omega_i \times ({}^i \omega_i \times {}^i P_{i+1}) + {}^i \dot{v}_i), \quad (6.47)$$

$${}^{i+1}\dot{v}_{C_{i+1}} = {}^{i+1}\dot{\omega}_{i+1} \times {}^{i+1}P_{C_{i+1}} + {}^{i+1}\omega_{i+1} \times ({}^{i+1}\omega_{i+1} \times {}^{i+1}P_{C_{i+1}}) + {}^{i+1}\dot{v}_{i+1}, \quad (6.48)$$

$${}^{i+1}F_{i+1} = m_{i+1} {}^{i+1}\dot{v}_{C_{i+1}}, \quad (6.49)$$

$${}^{i+1}N_{i+1} = {}^{C_{i+1}}I_{i+1} {}^{i+1}\dot{\omega}_{i+1} + {}^{i+1}\omega_{i+1} \times {}^{C_{i+1}}I_{i+1} {}^{i+1}\omega_{i+1}. \quad (6.50)$$

Inward iterations: $i : 6 \rightarrow 1$

$${}^i f_i = {}^{i+1}R {}^{i+1}f_{i+1} + {}^i F_i, \quad (6.51)$$

$${}^i n_i = {}^i N_i + {}^{i+1}R {}^{i+1}n_{i+1} + {}^i P_{C_i} \times {}^i F_i + {}^i P_{i+1} \times {}^{i+1}R {}^{i+1}f_{i+1}, \quad (6.52)$$

$$\tau_i = {}^i n_i^T {}^i \hat{Z}_i. \quad (6.53)$$



Manipulator Dynamics

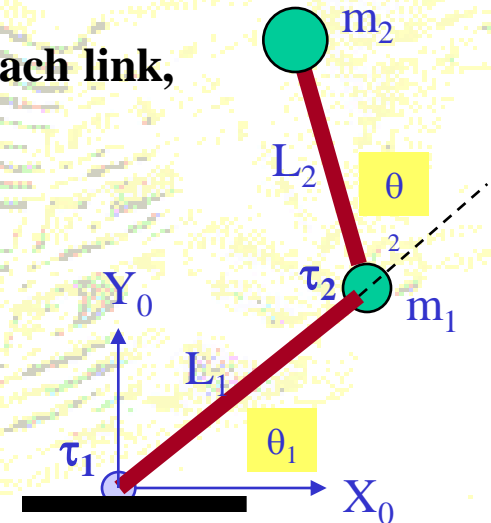
➤ Closed-form (Symbolic Form) Dynamic Equations:

Example: The 2-DOF Manipulator Arm.

– **Assumptions:** Point masses at the distal end of each link,

$${}^0\dot{v}_0 = g\hat{Y}_0 = \begin{bmatrix} 0 \\ g \\ 0 \end{bmatrix}, \quad (\text{gravity-term})$$

$$\begin{cases} {}^{c1}I_1 = 0 \\ {}^{c2}I_2 = 0 \end{cases} (\text{point-mass})$$



$$\begin{aligned} \tau_1 = & m_2 l_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) + m_2 l_1 l_2 C_2 (2\ddot{\theta}_1 + \ddot{\theta}_2) + (m_1 + m_2) l_1^2 \ddot{\theta}_1 - \\ & m_2 l_1 l_2 S_2 \dot{\theta}_2^2 - 2m_2 l_1 l_2 S_2 \dot{\theta}_1 \dot{\theta}_2 + m_2 l_2 g C_{12} + (m_1 + m_2) l_1 g C_1 \\ \tau_2 = & m_2 l_1 l_2 C_2 \ddot{\theta}_1 + m_2 l_1 l_2 S_2 \dot{\theta}_1^2 + m_2 l_2 g C_{12} + m_2 l_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) \end{aligned}$$

Actuator torques as a function of joints position, velocity, and acceleration.